1/6

<u>Magnetic Boundary</u> <u>Conditions</u>

Consider the **interface** between two different materials with dissimilar permeabilities:

 $\boldsymbol{H}_{\!\!1}(\overline{\boldsymbol{r}}), \boldsymbol{B}_{\!\!1}(\overline{\boldsymbol{r}})$

 $\boldsymbol{H}_{\!\scriptscriptstyle 2}(\overline{\boldsymbol{r}}), \boldsymbol{B}_{\!\scriptscriptstyle 2}(\overline{\boldsymbol{r}})$

 μ_{2}

 μ_1

Say that a magnetic field and a magnetic flux density is present in **both** regions.

Q: How are the fields in dielectric **region 1** (i.e., $H_1(\bar{r}), B_1(\bar{r})$) related to the fields in **region 2** (i.e., $H_2(\bar{r}), B_2(\bar{r})$)?

A: They must satisfy the magnetic boundary conditions !

First, let's write the fields at the interface in terms of their normal (e.g., $H_n(\bar{r})$) and tangential (e.g., $H_r(\bar{r})$) vector components:

$$H_{1n}(\bar{r}) \qquad H_{1}(\bar{r}) = H_{1r}(\bar{r}) + H_{1n}(\bar{r})$$

$$H_{1r}(\bar{r}) \qquad H_{1r}(\bar{r})$$

$$H_{2n}(\bar{r}) \qquad H_{2r}(\bar{r}) + H_{2n}(\bar{r})$$

 $\mu_{\rm 2}$

Our first boundary condition states that the **tangential** component of the magnetic field is **continuous** across a boundary. In other words:

$$\mathbf{H}_{1t}\left(\overline{\mathbf{r}_{b}}\right) = \mathbf{H}_{2t}\left(\overline{\mathbf{r}_{b}}\right)$$

where $\overline{r_b}$ denotes to **any** point along the interface (e.g., material boundary).

The tangential component of the magnetic field on one side of the material boundary is equal to the tangential component on the other side !

We can likewise consider the **magnetic flux densities** on the material interface in terms of their **normal** and **tangential** components:

$$\mathbf{B}_{1n}(\mathbf{\bar{r}}) \qquad \mathbf{B}_{1}(\mathbf{\bar{r}}) = \mu_{1}\mathbf{H}_{1}(\mathbf{\bar{r}})$$

$$\mu_{1} \qquad \mathbf{B}_{1r}(\mathbf{\bar{r}})$$

$$\mathbf{B}_{2n}(\mathbf{\bar{r}}) \qquad \mathbf{B}_{2r}(\mathbf{\bar{r}})$$

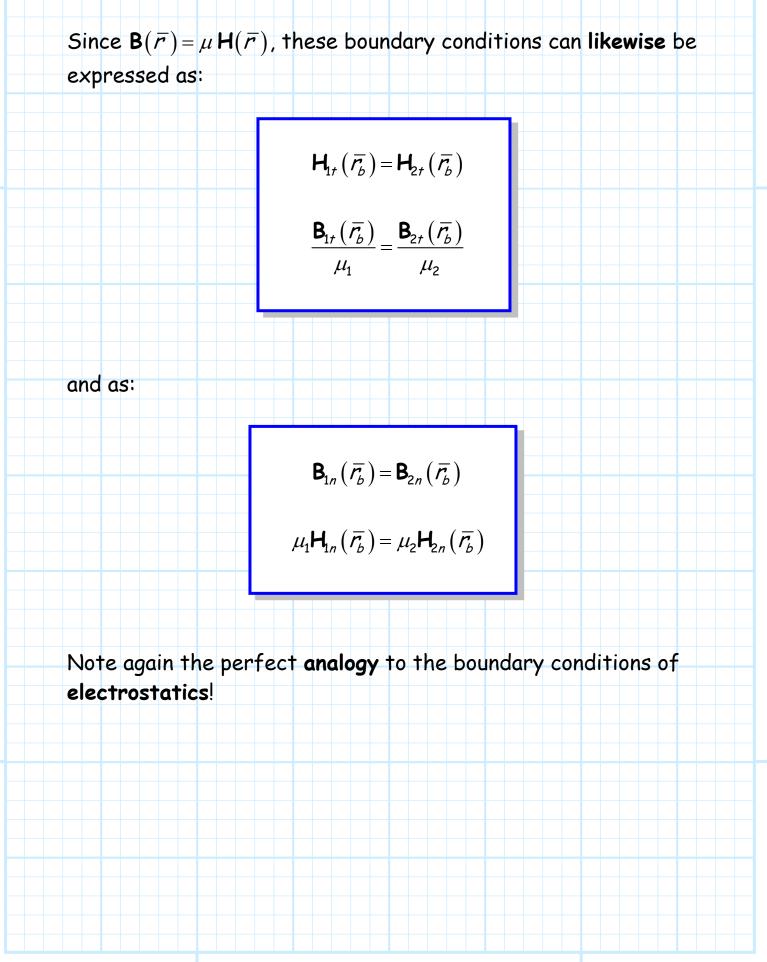
$$\mathbf{B}_{2n}(\mathbf{\bar{r}}) \qquad \mathbf{B}_{2}(\mathbf{\bar{r}}) = \mu_{2}\mathbf{H}_{2}(\mathbf{\bar{r}})$$

$$\mu_{2}$$

The second magnetic boundary condition states that the **normal** vector component of the **magnetic flux density** is **continuous** across the material boundary. In other words:

$$\mathbf{B}_{1n}\left(\overline{\mathbf{r}_{b}}\right) = \mathbf{B}_{2n}\left(\overline{\mathbf{r}_{b}}\right)$$

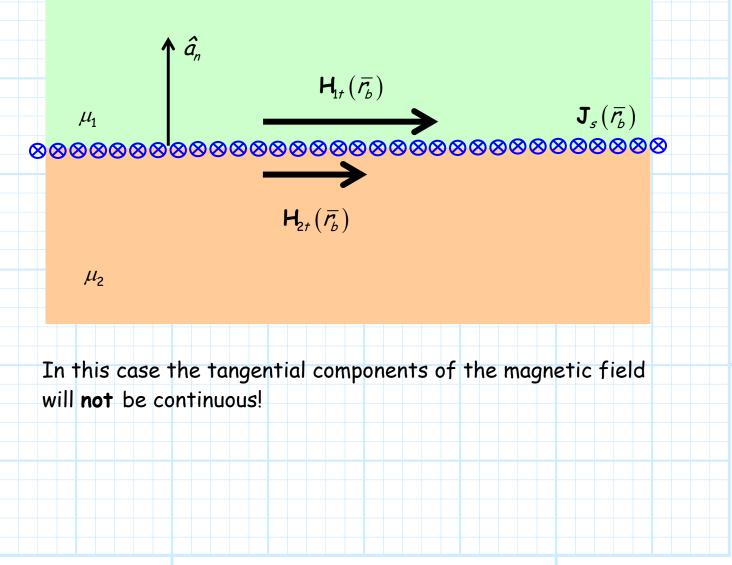
where $\overline{r_b}$ denotes **any** point along the interface (i.e., the material boundary).



Finally, recall that if a layer of **free charge** were lying at a dielectric boundary, the boundary condition for electric flux density was **modified** such that:

$$\hat{a}_{n} \cdot \left[\mathbf{D}_{1}(\overline{r_{b}}) - \mathbf{D}_{2}(\overline{r_{b}}) \right] = \rho_{s}(\overline{r_{b}})$$
$$\mathcal{D}_{1n}(\overline{r_{b}}) - \mathcal{D}_{2n}(\overline{r_{b}}) = \rho_{s}(\overline{r_{b}})$$

There is an **analogous** problem in magnetostatics, wherein a **surface current** is flowing at the interface of two magnetic materials:



Instead, they are related by the boundary condition:

$$\hat{a}_{n} \times \left(\mathsf{H}_{1}\left(\overline{r_{b}}\right) - \mathsf{H}_{2}\left(\overline{r_{b}}\right) \right) = \mathbf{J}_{s}\left(\overline{r_{b}}\right)$$

This expression means that:

1) $H_{lt}(\overline{r_b})$ and $H_{2t}(\overline{r_b})$ point in the same direction.

2) $H_{lt}(\overline{r_b})$ and $H_{2t}(\overline{r_b})$ are orthogonal to $J_s(\overline{r_b})$.

3) The difference between $|\mathbf{H}_{1t}(\overline{r_b})|$ and $|\mathbf{H}_{2t}(\overline{r_b})|$ is $|\mathbf{J}_s(\overline{r_b})|$.

Recall that $H(\bar{r})$ and $J_s(\bar{r})$ have the same units— Amperes/meter!

Note for this case, the boundary condition for the magnetic flux density remains **unchanged**, i.e.:

$$\mathsf{B}_{1n}\left(\overline{r_{b}}\right) = \mathsf{B}_{2n}\left(\overline{r_{b}}\right)$$

regardless of $\mathbf{J}_{s}(\overline{r_{b}})$.